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As further chemical and molecular structural studies of substituted carboranes become available, the nature of electronic and steric effects may become clearer in substitution reactions on the carboranes. In particular, models for activated complexes may be guessed, and the relevance of ground-state electronic structure gauged. This study indicates clearly that Cl substitution did not occur on the two most positively charged B atoms which were identified in both the molecular orbital and resonance studies.<sup>3,4</sup> Further studies are in progress. Acknowledgment.—The Fourier maps were calculated using Sly, Shoemaker, and Van den Hende's ERFR2 program, the least-squares refinements using G. S. Pawley's rigid body least-squares program, which is a special case of a more general treatment of rigid body motion in a crystal developed independently by V. Schomaker. We acknowledge support of this research by the Office of Naval Research, the National Science Foundation, and the U. S. Army Research Office (Durham). We wish to thank Dr. H. Schroeder for providing us with the sample.

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# Molecular Structure of Carboranes. A 1,2-Dicarbaclovododecaborane Derivative, $B_{10}H_{10}(CCH_2Br)_2$

BY DONALD VOET AND WILLIAM N. LIPSCOMB

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A three-dimensional X-ray diffraction study of single crystals of 1,2-bis(bromomethyl)carborane has proved the icosahedral structure of the  $B_{10}C_2$  unit. The covalently linked ethylenedecaborane  $B_{10}C_2$  unit is not correct, and the C-C distance of 1.64 Å, within the  $B_{10}C_2$  icosahedron is not abnormally short. The space group is  $P2_12_12_1$ , and there are four molecules in a unit cell having dimensions of a = 9.46, b = 10.12, and c = 13.78 Å.

The carboranes,<sup>1,2</sup> B<sub>10</sub>C<sub>2</sub>H<sub>12</sub>, their isomers,<sup>3-5</sup> and their derivatives form an unusually thermally stable group of compounds. The probable isoelectronic relation to the regular icosahedral  $B_{12}H_{12}^{-2}$  ion<sup>6</sup> led to the detailed prediction<sup>7-9</sup> of a near-icosahedral geometrical and valence structure, in which the C atoms of the  $B_{10}C_2$  unit participate in the electron deficiency of the assumed near-icosahedral framework. A detailed X-ray diffraction study became desirable when a distorted icosahedral structure was proposed<sup>3</sup> in which a short C-C bond about 1.54 Å. in length was assumed to distinguish one isomer, carborane, from a regular icosahedral isomer, neocarborane. Then, a detailed study became urgent when Zakharkin, et al., 10 published the results of an X-ray diffraction study which indicated that the  $B_{10}C_2$  unit of  $B_{10}H_{10}(CCH_2Br)_2$  had not a nearicosahedral structure, but an ethylenedecaborane covalent unit in which the two C atoms are joined to the outermost 6, 9 B atoms of the B<sub>10</sub> cage.

We show in the present study that the  $B_{10}C_2$  unit in

- (4) H. Schroeder and G. D. Vickers, *ibid.*, 2, 1317 (1963).
- (5) S. Papetti and T. L. Heying, J. Am. Chem. Soc., 86, 2295 (1964).
- (6) J. Wunderlich and W. N. Lipscomb, *ibid.*, **82**, 4427 (1960).
  (7) R. Hoffmann and W. N. Lipscomb, J. Chem. Phys., **86**, 3488 (1962).
- (8) R. Hoffmann and W. N. Lipscomb, Inorg. Chem., 2, 231 (1963).

(9) W. N. Lipscomb, "Boron Hydrides," The W. A. Benjamin Co., New York, N. Y., 1963.

(10) L. I. Zakharkin, V. I. Stanko, V. A. Brattson, Y. A. Chapovsky, and Y. T. Struchkov, Izv. Akad. Nauk SSSR, Ser. Khim., 11, 2069 (1963).  $B_{10}H_{10}(CCH_2Br)_2$  has a near-icosahedral structure in which the C–C distance of 1.64 Å. indicates not a single bond, but an electron-deficient type of bonding similar to that of the B atoms.

#### Structure Determination

1,2-Bis(bromomethyl)carborane was prepared by the method of Heying, *et al.*,<sup>1</sup> and then purified by vacuum sublimation at 100°. The melting point of 65– 66.5° is comparable with published values of  $^2$  68–69.5° and of  $^{10}$  66°. Also the infrared absorption spectrum of our sample was identical with that of a sample of the same material kindly supplied to us by Dr. H. Schroeder.

Single crystals were grown by sublimation at 60° onto a glass surface at room temperature. Two nearly spherical crystals, each about 0.2 mm. in diameter, were sealed into capillaries and aligned by X-ray diffraction methods along the a and c axes, respectively. These methods also established the reciprocal lattice symmetry as  $D_{2h}$ , the space group of  $P2_12_12_1$ , and the unit cell dimensions of  $a = 9.46 \pm 0.02$ ,  $b = 10.12 \pm 0.02$ , and  $c = 13.78 \pm 0.05$  Å.: systematic extinctions of h00 when h is odd, of 0k0 when k is odd, and of 00l when *l* is odd were noted on the zero level Weissenberg photographs, which were calibrated with an Al powder diffraction pattern. The reasonable assumption of four molecules per unit cell leads to the calculated density of 1.64 g. cm.<sup>-3</sup>, in essential agreement with a very rough value of 1.5 g. cm.<sup>-3</sup> found by flotation methods.

<sup>(1)</sup> T. L. Heying, J. W. Ager, Jr., S. L. Clark, D. J. Mangold, H. L. Goldstein, M. Hillman, R. J. Polak, and J. S. Szymanski, *Inorg. Chem.*, **2**, 1089 (1963).

<sup>(2)</sup> M. M. Fein, J. Bobinski, N. Mayes, N. Schwartz, and M. S. Cohen, *ibid.*, **2**, 1111 (1963).

<sup>(3)</sup> D. Grafstein and J. Dvorak, *ibid.*, 2, 1128 (1963).

The relative intensities of 1423 independent diffraction maxima were measured on levels Hkl for  $0 \leq H$  $\leq$  9 and *hkL* for  $0 \leq L \leq$  13 on the Buerger automated X-ray diffractometer at a scan rate<sup>11</sup> of  $1^{\circ}/\text{min.}$  and with the use of a Xe proportional counter. Background corrections, and then Lorentz-polarization corrections,  $(LP)^{-1} = 2 \cos^2 \mu \sin \Upsilon / (1 + \cos^2 2\theta)$ , were made in order to obtain the observed  $F_{hkl}$  (Table I). All final

#### TABLE I

#### LIST OF OBSERVED $F_{hkl}$ VALUES<sup>a</sup>

 $\begin{array}{c} \text{FABLE 1} \\ \textbf{LIST OF OBSERVED } F_{hkl} \text{ VALUES}^{a} \\ \hline \\ \begin{array}{c} \text{Fo} (K_1(1)_1(0)^{10}, 0)^{10}, 16^{10}, 93^{10}, 20^{10}, 10$  $\begin{array}{l} (7)(0,11), (15), (15), (15), (15), (11), (15), (12), (10), (12), (10), (15), (11), (12), (11),$ 4,10,7,6+( 5)(0, 4),9,\*,6,\*,10 H=11 (K)(L)+( 0)(0+4)+U+4,5+8+4+( 1)(0+4)+10+3,\*,\*,12+( 2)(0, 3),6,4+7+10+ ( 3)(0, 2),19,4+8

" The value of k and the range of l are given in parentheses. Unobserved reflections are indicated by an asterisk, and space group extinctions by an extra comma. All reflections are on the same scale as  $F_{000}$ , which is given as a reference point. No extinction or absorption corrections were made.

parameters were obtained from data correlated to a single scale, but the structure was solved from the a axis data which were assumed (justifiably, we later established) to be correlated by the instrument.

Solution of the structure presented no difficulties. The Br atoms were readily found in the sharpened threedimensional Patterson function, and a three-dimensional electron density map based upon phases of these Br atoms yielded fourteen peaks in the range of 210 to 329

(11) (a) J. A. Potenza and W. N. Lipscomb, J. Am. Chem. Soc., 86, 1874 (1964); (b) Inorg. Chem., 3, 1673 (1964).

units on an arbitrary scale on which Br atoms were 4668 and 4384. The highest false peak was 154. Thus the complete arrrangement of all B and C atoms was found without chemical assumptions, and it was immediately clear that twelve of these atoms were at the corners of a very nearly regular icosahedron.

We did not solve the absolute configuration of the crystal, but only record here that the positions of the two symmetry unrelated Br atoms, as found from the Patterson function, have a sixteenfold ambiguity associated with the arbitrary eightfold choice of the origin of the unit cell and the twofold choice of the enantiomorph. Thus our choice of enantiomorph is arbitrary.

Refinement of the structure proceeded rapidly. Values of  $R_F = \Sigma ||F_o|| - |F_c|| / \Sigma |F_o|$  for the 1353 *a* axis reflections were 0.35 for Br atoms only and 0.29when the fourteen peaks were initially given scattering amplitudes of B atoms. A new three-dimensional electron density map yielded the expected C atoms at heights of 445 to 520 and B atoms between 383 and 448 in height. Two cycles were carried out of least-squares refinement of the single scale factor, the 48 atomic position parameters, and the six translational, six rotational, and three center of rotational thermal parameters.12 Then separate scale factors were assigned to each level of reciprocal lattice data about the a axis, and two more cycles of three-dimensional least-squares refinement reduced  $R_F$  to 0.165. No further improvement was obtained by assigning separate anisotropic thermal parameters to each Br atom.

The data about the c axis became available, and a single scale of all reflections was then established. The 1239 reflections having  $|F_o| > 4$  and  $|F_c| > 4$  were used, and the 47 reflections having  $|F_o| < 4$  and  $|F_e| > 4$  were assigned  $F_{o} = 2$  in the final refinements, which yielded R = 0.132 after three cycles (Table II). All parameter

TABLE II					
Final Values for $R_F = (\Sigma    F_o  -  F_c   ) /  \Sigma   F_o )$					
		Range of			
hkl	R	$\sin \theta$	R		
A11	0.13	0.00 to 0.40	0.09		
h even	0.14	0.40 to 0.50	0.07		
h odd	0.12	0.50 to 0.60	0.09		
k even	0.13	0.60 to 0.65	0.12		
k odd	0.13	0.65 to 0.70	0.11		
<i>l</i> even	0.13	0.70 to 0.75	0.14		
<i>l</i> odd	0.13	0.75 to 0.80	0.20		
k + l even	0.13	0.80 to 0.85	0.28		
k + l odd	0.13	0.85 to 0.90	0.29		
h + l even	0.13	0.90 to 1.00	0.43		
h + l odd	0.13				
h + k even	0.13				
h + k  odd	0.13				
h + k + l even	0.13				
h + k + l odd	0.13				

shifts at this stage were less than  $0.7\sigma$ , where  $\sigma$  is the standard deviation. Weighting factors were w = 1 for  $|F_{\rm o}|^2 \leq 300$  and  $w = |F_{\rm o}|^2/300$  for  $|F_{\rm o}|^2 \geq 300$  where  $\sqrt{w}$  is the weighting factor in the equations of condition of the least-squares treatment.

(12) G. S. Pawley, Acta Cryst., 17, 457 (1964).

Final Atomic Coordinates in Fractions of Cell Lengths and Thermal Parameters <sup>a</sup>									
Atom	x	У	z	$10^{4}\beta_{11}$	$10^{4}\beta_{22}$	104 <i>β</i> 33	$10^{4}\beta_{12}$	$10^{4}eta_{13}$	$10^{4}\beta_{23}$
1	0.265	0.569	0.839	108	95	33	7	12	5
2	0.360	0.457	0.894	96	97	28	-10	7	-1
3	0.208	0.522	0.949	124	109	33	-10	20	13
4	0.088	0.554	0.857	105	123	50		13	21
5	0.172	0.506	0.745	119	121	36	-8	-3	22
6	0.344	0.450	0.770	100	96	27	-5	12	$^{2}$
7	0.245	0.342	0.950	123	106	33	-1	25	6
8	0.076	0.403	0.922	105	126	61	- 5	35	$^{2}$
9	0.053	0.401	0.793	90	141	68	-19		4
10	0.211	0.339	0.744	103	114	38	17	-5	$^{2}$
11	0.324	0.304	0.837	90	92	28	-7	10	-3
12	0.151	0.267	0.850	97	109	59	-11	19	-16
13	0.336	0.706	0.823	157	92	57	6	24	1
14	0.493	0.482	0.947	119	126	47	-25	-13	1
15	0.305	0.834	0.929	223	101	81	-29	19	4
16	0.665	0.476	0.871	93	153	97	-29	-3	-17

TABLE III

<sup>a</sup> The thermal parameters are in the form

 $\exp\left[-(h^2\beta_{11} + k^2\beta_{22} + l^2\beta_{33} + 2hk\beta_{12} + 2hl\beta_{13} + 2kl\beta_{23})\right]$ 

### **Results and Discussion**

Final parameters are given in Table III, distances in Table IV, and angles in Tables V and VI. The atomic numbering system is shown in Fig. 1. Corrections for

TABLE IV					
BOND DISTANCES <sup>4</sup>					
Bond	Corrected length, <sup>b</sup> Å.	Bond	Corrected length, <sup>b</sup> Å.		
	C-C	B-	·B		
12	1.64	3-4	1.74		
1–13	1.55	3-7	1.86		
2 - 14	1.49	3–8	1.78		
	CBr	4-5	1.80		
13 - 15	1.96	4-8	1.77		
14–16	1.94	4-9	1.81		
	B-C	5-6	1.76		
1 - 3	1.68	5-9	1.69		
1-4	1.70	5 - 10	1.74		
1 - 5	1.69	6 - 10	1.73		
1 - 6	1.71	6 - 11	1.76		
2 - 3	1.75	7-8	1.76		
2-6	1.71	7-11	1.77		
2-7	1.77	7-12	1.80		
2 - 11	1.77	8-9	1.78		
		8-12	1.84		
		9 - 10	1.75		
		9-12	1.81		
		10 - 11	1.71		
		10 - 12	1.72		
		11 - 12	1.69		

<sup>a</sup> Standard deviations (Å.) for the various types of bonds are: B-C,  $\sigma = 0.037$ ; B-B,  $\sigma = 0.033$ ; B-C(methyl),  $\sigma = 0.045$ ; and C(methyl)-Br,  $\sigma = 0.016$ . The standard deviations were computed by averaging all bond lengths of the same type and applying the equation

$$\sigma = \left(\frac{\sum_{i=1}^{N} (x_i - \bar{x})^2}{N - 1}\right)^{1/2}$$

<sup>b</sup> Corrections for thermal vibrations were computed using the least-squares rotation tensor and were found to be small compared to the standard deviation in bond lengths. For B-B, B-C, B-C(methyl), and C(methyl)-Br bonds, the largest corrections were 0.0044, 0.0045, 0.0013, and 0.0017 Å., respectively.



Fig. 1.-The molecular structure and numbering scheme for  $B_{10}H_{10}(CCH_2Br)_2$ . B atoms are designated as single circles, C atoms are filled circles, and Br atoms are double circles.



Fig. 2.—The molecular packing of  $B_{10}H_{10}(CCH_2Br)_2$  as seen in the projection along the a axis.

molecular thermal oscillations are only about one-tenth of the standard deviations of the bond distances (footnotes, Table IV). Bond angles are listed in a form suitable for construction of a model.

Our major conclusion is that both bond distances and

TABLE V

	Model Dr	ULLING ANGLES	
Center	Coordi-	4 a	a
atom	atom	φ, deg.	ρ, deg.
1	2	90 0	0.0
1	3	0.0	64 0
	4	322.5	114.9
	7 5	959 9	114.2
	6	202.2	61 5
0	1	214.4	01.0
4	1	90.0	0.0 50.1
	3 C	145 0	09.1 01 F
	0	145.6	6.10
	(	41.4	108.6
0	8	41.1	70.6
3	1	90.0	0.0
	2	0.0	56.9
	4	219.0	59.8
	8	262.5	103.2
	9	259.3	67.1
4	1	90.0	0.0
	3	0.0	58.2
	5	218.1	57.6
	9	255.1	100.9
	10	255.1	66. <b>5</b>
5	1	90.0	0.0
	4	0.0	58.4
	6	215.3	59.5
	10	254.4	105.6
	11	254.6	70.5
6	1	90.0	0.0
	2	0.0	57.1
	5	140.0	58.3
	7	37.4	69.1
	11	40.2	104.9
7	2	90.0	0.0
	3	0.0	57.8
	8	38.9	103.5
	11	139.5	59.9
	12	103.2	104.0
8	3	90.0	0.0
	4	0.0	58.7
	7	140.7	63.2
	9	38.2	107.6
	12	101.8	110.1
9	4	90.0	0.0
	5	0.0	61.8
	8	136.1	59.1
	10	38.0	110.0
	12	99.6	109.8
10	5	90.0	0.0
	6	0.0	60.8
	9	139.8	57.9
	11	39.4	109.1
	12	102.3	109.5
11	<b>2</b>	90.0	0.0
	6	0.0	58.0
	7	217.2	60.0
	10	321.8	105.6
	12	257.8	109.2
12	7	90.0	0.0
	8	0.0	57.6
	9	35.1	105.4
	10	97.5	108.2
10	11	136.1	61.0
13	1	90.0	0.0
4.4	15	0.0	114.7
14	2	90.0	0.0
	6	0.0	115.8

<sup>a</sup> See ref. 4, p. 228, for definitions of  $\phi$  and  $\rho$ . These are spherical coordinates in which  $\phi$  is measured in the *x*, *y* plane, and  $\rho$  is measured from the *z* axis.

TABLE VI

Average Bond Angles (deg.) <sup>4</sup>					
Bond type	A verage angle	Largest angle	Smallest angle		
B-B-B	60	63	56		
B-B-C	59	61	58		
B-C-B	62	64	60		
B-C-C	62	64	59		
C-B-C	57	57	57		
B-C-C(methyl)	119	125	111		
C-C-C(methyl)	121	125	117		
C-C(methyl)-Br	115	116	115		
	AVERAGE BOND Bond type B-B-B B-C-B B-C-B B-C-C C-B-C B-C-C(methyl) C-C-C(methyl)-Br	AVERAGE BOND ANGLES (DE           Bond         Average           type         angle           B-B-B         60           B-B-C         59           B-C-B         62           B-C-C         62           C-B-C         57           B-C-C(methyl)         119           C-C-C(methyl)         121           C-C(methyl)-Br         115	AVERAGE BOND ANGLES (DEG.) <sup>4</sup> Bond       Average angle       Largest angle         B-B-B       60       63         B-B-C       59       61         B-C-B       62       64         B-C-C       57       57         B-C-C(methyl)       119       125         C-C(methyl)-Br       115       116		

<sup>*a*</sup> The bond angles were computed using values for the atomic positions which were corrected for thermal vibrations by means of the least-squares rotation tensor. These corrections were found to be small compared to the standard deviation of the bond angles, the largest correction being  $0.10^{\circ}$ . Estimated standard deviations of bond angles are about  $3^{\circ}$ .

angles clearly indicate the near-icosahedral geometry of the  $B_{10}C_2$  unit, not the ethylenedecarborane structural unit (Fig. 2 of ref. 11b derived in the X-ray study reported by Zakharkin, *et al.*<sup>10</sup>).

Secondly, our results do not support the proposal<sup>3</sup> that the carborane cage has a 1.54 Å. C-C distance, which opens up to the regular icosahedral distance of 1.77 Å. in the transformation at about 470° to neocarborane. The C–C distances of  $1.64 \pm 0.40$  Å. in this study, and of  $1.67 \pm 0.02$  Å. in octachlorocarborane,<sup>11</sup> leave no physical basis for a high barrier between these proposed structures, and hence the proposal<sup>13</sup> of a rearrangement through a near-cuboctahedral intermediate (perhaps involving localized regions of expansion and contraction) seems most reasonable, at present, for the relation between carborane and neocarborane. An X-ray diffraction study of a neocarborane structure is under way. With almost improper respect for the observation that predictions in science are more difficult than in science fiction (which is limited to what we can imagine), we maintain<sup>9,13</sup> our present belief that neocarborane has its two C atoms in the meta positions, neither adjacent nor opposite in the icosahedron.

The crystal structure (Fig. 2) has only one feature of interest: a  $Br_{15} \cdots Br_{16}$  contact of 3.72 Å. between different molecules. This is the only intermolecular distance which is less than the expected van der Waals contact, which is 3.9 Å. for Br. Br. There are many similar examples in other crystal structures containing halogen atoms, including the crystals of the diatomics  $Cl_2$ ,  $Br_2$ , and  $I_2$ , but not  $F_2$ . It may be that this particularly strong intermolecular interaction is responsible in part for the distortion of bond angles  $C_2C_1C_{13}$  =  $116.6^{\circ}$  and  $C_1C_2C_{14} = 125.0^{\circ}$  from each other and from the value of 121.7° expected from a regular icosahedral structure. The implied easy bending of external bonds suggests that the nondirectional character of electrondeficient bonding already indicated by the bond distances produces little restoring torque for angular deformations. Further evidence for this lack of directionality of bonding within the polyhedral surface is sug-

(13) R. Hoffmann and W. N. Lipscomb, J. Chem. Phys., 37, 2872 (1962).

#### Vol. 3, No. 12, December, 1964

gested by the rearrangements and by the great simplicity of molecular spectra found in these complex polyhedral molecular species.

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## Geometrical Theory of Boron Hydrides

#### By WILLIAM N. LIPSCOMB

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Based upon the assumptions that in some probable new boron hydrides the geometrical and bonding environments of B atoms will be similar to those in known hydrides, that all stable boron hydride molecules encounter external H...H contacts only upon collision with other molecules of the same or similar type, and finally that all intramolecular H...H contacts are reasonable ( $\geq$  about 2 Å.) it is shown here that only a very limited number of predicted new boron hydrides based upon a single icosahedral fragment can be expected. Geometrical structural predictions are given for these possible new boron hydrides.

The three-center resonance and molecular orbital theories<sup>1</sup> of boron hydrides take the steric properties of boron hydride structures into account in only a general and not very restrictive way. Accordingly, it seems appropriate to give specific attention to the problems of intramolecular and intermolecular steric requirements, within the framework of these valence theories, in the hope that only a limited number of predictions would result. This paper represents a first theory of this type and is restricted to neutral single icosahedral fragments for the boron arrangements. We have therefore specifically excluded boron hydride ions, fragments, or polyhedra joined by BBB or BB<sup>2</sup> bonds  $(e.g., B_{10}H_{16})$ , by single or double hydrogen bridges<sup>3</sup>  $(e.g., B_{20}H_{19}^{-})$ , by the sharing of two B atoms<sup>1</sup> (e.g., $B_{18}H_{22}$  and  $i-B_{18}H_{22}$ ) or of four B atoms (e.g.,  $B_{20}H_{16}$ ), or fragments obtained from possible degradation of structures in such a way that these joinings are preserved. Equivalent theories of these additional kinds of molecules and ions can easily be derived by similar methods.

Accordingly, the assumptions and procedures are as follows. (1) Construct models of the known boranes, ions, and ligand derivatives from balls precisely drilled according to the internuclear straight lines between all nearest neighboring atoms. These angles have been listed<sup>1,4</sup> for all of the known compounds, including<sup>5</sup> B<sub>8</sub>H<sub>12</sub>. (2) Choose from these models all different types of balls

(2) R. Grimes, F. E. Wang, R. Lewin, and W. N. Lipscomb, Proc. Natl. Acad. Sci. U. S., 47, 996 (1961).

(3) W. N. Lipscomb, ibid., 47, 1791 (1961).

(4) R. D. Dobrott, L. Friedman, and W. N. Lipscomb, J. Chem. Phys., 40, 866 (1964); see footnote 3.

(5) R. E. Enrione, F. P. Boer, and W. N. Lipscomb, J. Am. Chem. Soc., 86, 1451 (1964); Inorg. Chem., 3, 1659 (1964).

for support of this research. Also, we thank G. S. Pawley for discussions and for use of the rigid body least-squares refinement program which reduces thermal parameters to 15 per molecule, and which is a special case of a theory by V. Schomaker which includes, in addition, the six additional screw-rotational tensor components omitted in these refinements.

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#### TABLE I

#### NEAREST COORDINATION ENVIRONMENTS OF B ATOMS OF ALL TYPES в $H_t^a$ $H_b^a$ Example $\mathbf{2}$ 2 1 $B_{3}H_{f}^{-}$ , the symmetry related B atoms $\mathbf{2}$ $\mathbf{2}$ $\mathbf{2}$ $B_3H_8^-$ , the unique B 3 1 1 B<sub>6</sub>H<sub>10</sub>, nonapical atoms joined with the short bond 3 1 $\mathbf{2}$ $B_4H_{10}$ , the inner pair 3 $\mathbf{2}$ 0C<sub>2</sub>H<sub>5</sub>NH<sub>2</sub>B<sub>8</sub>H<sub>12</sub>NHC<sub>2</sub>H<sub>5</sub>, B joined to C<sub>2</sub>H<sub>5</sub>- $NH_2$ 3 $\mathbf{2}$ B<sub>9</sub>H<sub>13</sub>NCCH<sub>3</sub>, the BH<sub>2</sub> groups 1 4 1 0 B<sub>9</sub>H<sub>13</sub>NCCH<sub>3</sub>, the BH on the symmetry plane 4 1 1 B<sub>9</sub>H<sub>15</sub>, adjacent to the BH<sub>2</sub> group 4 $\mathbf{2}$ 0 B₅H11, apex $\mathbf{5}$ 1 0 B6H10, apex (5)0 $(0)^{b}$ B10H16, atoms joining the B6H8 units (50 $2)^{b}$ i-B<sub>18</sub>H<sub>22</sub>, on the molecular twofold axis (6 0 $(0)^{b}$ B (elementary), not yet found in the hydriđes (6 0 $1)^{b}$ B<sub>18</sub>H<sub>22</sub>, nearest molecular center of symmetry (7

0 0)6 i-B<sub>18</sub>H<sub>22</sub>, on the molecular twofold axis

<sup>a</sup> The subscripts t and b refer to terminal and bridge H atoms, respectively. <sup>b</sup> Atoms of types listed in parentheses are required for structure types not considered in this paper, e.g., structures based upon polyhedra or fragments joined by boron-boron bonds, by one or more bridge H atoms, or by common B atoms.

(Table I), but, for the single icosahedral frameworks considered here, limit the choices to B atoms having at least one terminal H atom  $(H_t)$ . (3) Construct all possible models with the use of B-B distances of about 1.8 Å.,  $B-H_t$  distances of 1.2 Å., and  $B-H_b$  (bridge) distances of about 1.4 Å., on a scale of 1 Å. = 10 cm. (4) Reject all models which have intramolecular H...H contacts of less than about 2.0 Å, where these H atoms

<sup>(1)</sup> W. N. Lipscomb, "Boron Hydrides," The W. A. Benjamin Co., New York, N. Y., 1963, p. 227 ff.